# Characterization of MoS<sub>2</sub>−Graphene Composites for High-Performance Coin Cell Supercapacitors

Mark A. Bissett,\*<sup>,†</sup> Ian A. Kinloch,<sup>‡</sup> and Robert A. W. Dryfe<sup>\*,†</sup>

† School of Chemist[ry,](#page-8-0) University of Manchester, Oxford Road, Manchester [M1](#page-8-0)3 9PL, United Kingdom ‡ School of Materials, University of Manchester, Oxford Road, Manchester M13 9PL, United Kingdom

**S** Supporting Information

[AB](#page-8-0)STRACT: [Two-dimensio](#page-8-0)nal materials, such as graphene and molybdenum disulfide  $(MoS<sub>2</sub>)$ , can greatly increase the performance of electrochemical energy storage devices because of the combination of high surface area and electrical conductivity. Here, we have investigated the performance of solution exfoliated  $MoS<sub>2</sub>$ thin flexible membranes as supercapacitor electrodes in a symmetrical coin cell arrangement using an aqueous electrolyte  $(Na_2SO_4)$ . By adding highly conductive graphene to form nanocomposite membranes, it was possible to increase the specific capacitance by reducing the resistivity of the electrode and altering the morphology of the membrane. With continued charge/discharge cycles the performance of the membranes was found to increase significantly (up to 800%), because of partial reexfoliation of the layered material with continued ion intercalation, as well as



increasing the specific capacitance through intercalation pseudocapacitance. These results demonstrate a simple and scalable application of layered 2D materials toward electrochemical energy storage.

KEYWORDS: graphene, MoS<sub>2</sub>, supercapacitors, pseudocapacitors, electrochemistry, energy storage, coin cell

# **ENTRODUCTION**

Two-dimensional (2D) materials such as graphene and molybdenum disulfide  $(MoS<sub>2</sub>)$  are excellent candidates for applications in electrochemical energy storage because of their high available surface area and versatile electronic structure. Porous carbon materials, and graphene in particular, have been investigated previously for applications in electrochemical energy storage because of their many beneficial properties such as mechanical strength, high specific surface area (2675  $(m^2/g)$ , and high thermal and electrical conductivity.<sup>1–6</sup> Other layered 2D materials such as transition metal dichalcogenides (TMDs) including  $MoS<sub>2</sub>$  have been investigated t[o a](#page-8-0) lesser extent but are gaining increased interest recently for integration into electronic devices due to similarities with graphene.<sup>7,8</sup> Exfoliated  $MoS<sub>2</sub>$  in particular possesses high catalytic activity which makes it an efficient hydrogen evolution catalyst $9-12$  [as](#page-8-0) well as a useful energy storage material for use in lithium and sodium ion batteries.<sup>13,14</sup> These layered 2D materials [c](#page-8-0)a[n](#page-9-0) be synthesized in several different ways including micromechanical peeling<sup>15</sup> and chemic[al va](#page-9-0)por deposition,<sup>15,16</sup> but currently the most applicable for cost-effective, large scale production is throug[h t](#page-9-0)op down exfoliation of bulk [powd](#page-9-0)ers via ultrasonic treatment in specially chosen solvents. This solution phase exfoliation has been demonstrated previously for graphene $17$ and can be extended to other 2D materials<sup>18,19</sup> to create large volumes of monolayer and few layer flakes that can then [be](#page-9-0) deposited onto a substrate or formed into fi[lms](#page-9-0).

Supercapacitors are one application for which 2D materials can greatly increase performance and efficiency. Supercapacitors are unique because they provide a much higher power density than conventional batteries while also having a higher energy density than regular capacitors while still maintaining almost limitless cycling durability. With the continued miniaturization of electronic devices the need to create similarly miniature power sources is increasing, and supercapacitors based on 2D materials fulfill these requirements by having high specific capacitances for small active areas and masses. The charge storage mechanism in supercapacitors can be divided into two groups; (1) electric double-layer capacitors (EDLCs) and (2) pseudocapacitors. EDLCs store energy through the accumulation of ions at the interface between the electrode and the electrolyte, and because of this, their performance is dependent on the surface area of the electrode, pore size, and electrical conductivity.<sup>20</sup> Pseudocapacitors, on the other hand, use fast reversible redox reactions at the electrode−electrolyte interface, resul[tin](#page-9-0)g in Faradaic charge transfer. EDLCs have the advantage of almost limitless cyclability with little degradation but comparatively low specific capacitances when compared to pseudocapacitors which have much higher specific capacitance but often suffer from performance degradation with repeated charge−discharge cycles.21−<sup>23</sup>

To enable comparison between different electrode materials, specifi[c capa](#page-9-0)citance values are typically normalized by either the mass of the active material  $(F/g)$ , the projected electrode area  $(mF/cm<sup>2</sup>)$ , or volume  $(F/cm<sup>3</sup>)$ . However, care should be taken

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Received: May 28, 2015
Accepted: July 21, 2015
Published: July 21, 2015
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when comparing specific capacitance values reported in literature, as very low masses  $(\mu g)$  can give greatly overestimated values when extrapolated to the gram or even kilogram scale, $6,24$  as well as excluding the contribution of conductive additives (such as carbon black) or binding agents. Because of th[e](#page-8-0) [di](#page-9-0)fficulties in scaling up specific capacitance values for very small masses ( $\mu$ g or mg) of active material to large scale (g or kg) applications, the use of area normalized capacitance  $(mF/cm^2)$  is regarded as more accurately scalable, $1,25,26$  particularly for thin films of electrode material. The use of three-electrode systems can also exaggerate the perform[a](#page-8-0)[nce o](#page-9-0)f an electrode, more than doubling the measured capacitance and not accurately reflecting the performance of real world devices.<sup>24,25</sup> Thus, to accurately predict the performance of an electrode material, testing should be undertaken in a sy[mme](#page-9-0)trical two electrode system without any additives or binders.

Despite the increasing number of studies on the use of  $MoS<sub>2</sub>$ in supercapacitor electrodes, very little is understood about the fundamental electrochemical processes that contribute to its electron transfer characteristics and pseudocapacitance.<sup>27,28</sup> As well as Faradaic electron transfer resulting in the change in oxidation states of the Mo and S atoms at edge sites, [seve](#page-9-0)ral different non-Faradaic charge storage mechanisms have been proposed, including the adsorption of electrolyte ions onto the  $MoS<sub>2</sub>$  surface as well as the intercalation of ions between neighboring layers.<sup>29−31</sup> Previous work on thin layer graphenebased supercapacitors has reported a range of specific capacitances from [2 to](#page-9-0) 5  $mF/cm<sup>2</sup>$  depending on the testing parameters, electrolyte, and electrode architecture.<sup>32-34</sup> However, there have been fewer published values for  $MoS<sub>2</sub>$  based devices but they range from 2 to 3.1  $mF/cm<sup>2</sup>$  [for so](#page-9-0)lution exfoliated MoS<sub>2</sub><sup>31,35</sup> and 8 mF/cm<sup>2</sup> for hydrothermally grown material. $31$  Recently, the use of the metallic metastable 1T- $MoS<sub>2</sub>$  phase wa[s dem](#page-9-0)onstrated to be an ideal supercapacitor electrod[e m](#page-9-0)aterial due to its high conductivity combined with ion intercalation pseudocapacitance.<sup>29</sup> With the obvious benefits of these 2D materials it is possible to combine graphene with  $MoS<sub>2</sub>$  to produce hig[hly](#page-9-0) efficient composite electrochemical supercapacitor devices. These composite devices are able to produce high gravimetric specific capacitances between 205 and 282 F/g.36−<sup>38</sup> Similar composites of graphene and metal oxides  $(e.g., MnO<sub>2</sub>/$ graphene) have also been demonstrated previo[usly, a](#page-9-0)nd these maximize the available pseudocapacitance of the oxide through the interconnectivity of both materials.39−<sup>41</sup>

In this work, we have produced few-layer (1–3 layers)  $MoS<sub>2</sub>$ by solution phase exfoliation; this w[as](#page-9-0) [the](#page-9-0)n combined with dispersions of exfoliated graphene to produce a  $MoS<sub>2</sub>/graphene$ composite electrode material. The exfoliated graphene and  $MoS<sub>2</sub>$  were first characterized by a variety of techniques including optical absorption, Raman spectroscopy, as well as Xray photoelectron spectroscopy (XPS) and scanning electron microscopy (SEM) to determine the quality of the electrode materials. Membranes of graphene and  $MoS<sub>2</sub>$  were formed by filtration of the dispersions and were characterized in symmetrical electrochemical coin cells with aqueous electrolytes to measure the charge storage behavior. Because of the inherent conductivity of the thin membranes, they can act as the active electrode without the need for any binding agents, simplifying device design while also reducing overall device weight. It was also possible to directly stack the membranes back-to-back, negating the need for a separate ion porous separator. It was found that while the individual materials had moderate specific capacitance values, the composite membranes had significantly improved values depending on the relative ratio of  $MoS<sub>2</sub>/graphene$ , combined with excellent cyclability. Because of repeated ion intercalation and deintercalation that occurs with repeated cycling, the performance of the membranes increased significantly (up to 800%) with cycling, attributed to further exfoliation of the layered materials. The improved performance of the composite electrodes is attributed to the decreased resistance from the incorporated graphene flakes combined with the increased surface area created by the stacking interactions of the  $MoS<sub>2</sub>$  and graphene flakes. These results indicate that composites of 2D materials can be integrated into highly efficient energy storage applications, which with further optimization should be suitable for widespread applications.

#### **EXPERIMENTAL SECTION**

 $MoS<sub>2</sub>$  Dispersion Preparation.  $MoS<sub>2</sub>$  dispersions were created by liquid phase exfoliation as has been reported previously.<sup>18,19,42</sup> Briefly,  $MoS<sub>2</sub> powder (10 mg/mL) was dispersed in N-methyl-2-pyrrolidone$ (NMP) and ultrasonically processed in a bath sonicat[or \(Elm](#page-9-0)asonic P70H) operating at 37 kHz and 40% amplitude for 12 h while cooling to maintain a stable temperature of 20 °C. After this sonication, the dispersions were centrifuged at 6000 rpm (3139g) for 30 min to remove any unexfoliated material, the supernatant was then decanted and fresh solvent added before repeating the centrifugation to ensure a narrow distribution of flake dimensions and thickness. The resultant  $MoS<sub>2</sub>$  suspensions were stable in solution for several months with no detectable sedimentation.

Graphene Dispersion Preparation. Graphene powder (Grat-G1L, BGT Materials, U.K.) was dispersed in NMP (10 mg/mL) before sonicating briefly (20 min) to ensure a homogeneous suspension. This graphene dispersion was then centrifuged twice at 6000 rpm (3139g) for 30 min, as performed for the  $MoS<sub>2</sub>$ , to ensure that no unexfoliated material remained. The resultant graphene suspensions were also stable in solution for several months with almost no detectable sedimentation.

Raman Spectroscopy and Photoluminescence. To ascertain the  $MoS<sub>2</sub>$  and graphene layer number, the dispersions were deposited onto a  $Si/SiO<sub>2</sub>$  (300 nm thickness of oxide) wafer and the solvent was removed by heating on a hot plate allowing flakes to be identified by optical microscopy and Raman spectroscopy. Confocal Raman spectroscopy and photoluminescence measurements were conducted using a Renishaw inVia microscope with a 532 nm (2.33 eV) excitation at a power of 1 mW with a 100× objective and a grating of 1800 l/mm to achieve a spectral resolution of  $\sim$ 1 cm<sup>-1</sup>. .

**Membrane Preparation.** Films of the exfoliated  $MoS<sub>2</sub>$  and graphene, as well as mixtures of both, were synthesized by first diluting the NMP dispersions in isopropanol (IPA) by a factor of 20 followed by filtering through preweighed PVDF filters with 0.1  $\mu$ m pore size. These were then washed with IPA to remove any NMP residue and dried in an oven at 80 °C to remove any residual solvent. The mass of active material used was determined by weighing the membrane after drying and kept at ~0.75 mg (mass loading: 1 mg/cm<sup>2</sup>). The MoS<sub>2</sub>, graphene, and composite films were then each analyzed by Raman spectroscopy, XPS, SEM, and electrochemistry.

Electrochemistry. Cyclic voltammetry (CV), electrochemical impedance spectroscopy (EIS), and galvanostatic charge/discharge (GCD) were performed using a PGSTAT302N potentiostat (Metrohm Autolab, The Netherlands). All electrochemical measurements were performed using a sealed symmetrical coin cell (CR2032). The membranes were stacked back-to-back within the coin cell with the active material making direct contact with the current collector. The electrolyte used was aqueous 1 M  $Na<sub>2</sub>SO<sub>4</sub>$ . Cyclic voltammetry was performed at scan rates ranging from 5 to 1000 mV/s. EIS was performed at a frequency range of 0.1 Hz to 100 kHz with a 10 mV (rms) perturbation and 0 V dc bias.



Figure 1. (a) Illustration showing the crystal structure of  $MoS<sub>2</sub>$  and the sonication process along with a photograph showing the stable  $MoS<sub>2</sub>$ dispersion in NMP after centrifugation. (b) Optical absorption spectrum of MoS<sub>2</sub> dispersion in NMP showing the A and B excitonic absorptions. (c) Raman spectrum of solution exfoliated MoS2 flake supported on SiO2. The main Raman peaks are labeled and the vibrations shown schematically.

Scanning Electron Microscopy. Scanning electron microscope (SEM) analysis was performed on a FEI/Philips XL30 E-SEM under high vacuum conditions with an accelerating voltage of 15 kV. All images were taken using the secondary electron detector.

X-ray Photoelectron Spectroscopy. X-ray photoelectron spectroscopy (XPS) was performed at the EPSRC NEXUS facility (Newcastle, U.K.) using an AXIS Nova spectrometer (Kratos Analytical, U.K.) which utilizes a monochromatic Al K $\alpha$  source operated at 225W (15 kV and 15 mA) with an analysis area of 300  $\mu$ m  $\times$  700  $\mu$ m, and three points were analyzed on each sample. Survey spectra were collected at a pass energy of 160 eV and were the average of three sweeps, while high resolution spectra were collected at a pass energy of 20 eV and were the average of 10 sweeps. Charge neutralization was used for all samples.

### ■ RESULTS AND DISCUSSION

Synthesis and Characterization of Electrode Membranes. To understand the performance of any energy storage devices based on  $MoS_2$  flakes, it is paramount to first thoroughly characterize the active material. As shown in the schematic of Figure 1a, dispersions of  $MoS<sub>2</sub>$  were prepared by solution phase exfoliation of bulk  $MoS<sub>2</sub>$  powder by ultrasonication in N-methyl-2-pyrrolidone (NMP). Because of the relatively weak interlayer van der Waals interactions binding the  $MoS<sub>2</sub>$  layers, exfoliation down to a few layers or monolayers by ultrasonication in the correctly chosen solvent is possible. NMP was chosen, as it has been shown to be highly efficient at exfoliating  $MoS<sub>2</sub>$  as well as maintaining a stable disper $s$ ion.<sup>18,19,42</sup> As the bulk powder is reduced in thickness down to a few layers (1−3), it is stabilized in solution and prevented fro[m reagg](#page-9-0)regation due to the close match in surface energy between the  $MoS<sub>2</sub>$  layers and the solvent. This close match in surface energy between the  $MoS<sub>2</sub>$  flakes and solvent causes the enthalpy of mixing to be small, ensuring a stable dispersion.<sup>18,19,42</sup> The optical image in Figure 1a shows the stable  $MoS<sub>2</sub>$  dispersion in NMP with no discernible sedimentation over ex[tended](#page-9-0) periods of time. Figure 1b shows the absorption spectrum of the  $MoS<sub>2</sub>$  dispersion with two clear excitonic absorptions at 1.85 and 2 eV labeled A and B, respectively. The lack of any convolution of these two excitonic absorptions indicates a lack of damage or disorder that may have occurred because of the ultrasonication process.<sup>43</sup> This indicates that

solution exfoliation does not significantly damage the crystal structure.

Figure 1c shows the Raman spectrum of an exfoliated  $MoS<sub>2</sub>$ flake after a drop of the dispersion was drop cast, after further dilution, onto a silicon wafer  $(Si/SiO<sub>2</sub>)$  to allow easy identification by optical microscopy. The two characteristic Raman peaks for  $\mathrm{MoS}_{2}$  are labeled the in-plane vibration  $(\mathrm{E_{2g}^{1}})$ and the out-of-plane  $(A_{1g})$ .<sup>44,45</sup> The wavenumber of these</sup> peaks, as well as peak separation, can be used to determine the number of  $MoS<sub>2</sub>$  layers. T[ypica](#page-9-0)lly the Raman spectrum of micromechanically exfoliated  $MoS<sub>2</sub>$  crystals exhibits a peak separation of  $\sim$ 20 cm<sup>-1</sup> while bulk MoS<sub>2</sub> has a separation of  $\sim$ 25 cm $^{-1}$ , as well as the full width at half-maximum (fwhm) of the peaks increasing with decreasing layer number. $44,45$  When we compare the Raman spectrum of the raw powder and the exfoliated flakes of  $MoS<sub>2</sub>$  (Supporting Information, [Figu](#page-9-0)re S1), we see a significant shift due to the sonication process indicating that the resultant fl[akes are few layer \(](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)1−3 layers). The flake size and thickness are dependent on both the sonication time used and also the postprocessing such as centrifugation speed and duration.<sup>42</sup> Thus, these parameters can be easily tuned to produce films with the appropriate dimensions. We also see a signifi[ca](#page-9-0)nt difference when we compare the spectrum of the bulk crystal compared to the raw powder (Figure S1), where the peak frequencies are downshifted for the powder as well as significantly broader. This indicates [that the p](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)owder is much more disordered than the bulk material, as may be expected. By varying the crystallinity of the starting material, it may be possible to reduce the distribution of the lateral flake size and thickness as needed for a specific application. Because of the semiconducting nature of  $MoS<sub>2</sub>$ , we can also measure the photoluminescence (PL) of the bulk starting materials and compare this to the exfoliated material as seen in the Supporting Information (Figure S2) which shows a clear shift in electronic structure due to the reduction in thickness. B[ecause of the large numbe](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)rs of flakes in any given dispersion, there exists a distribution of thicknesses and also lateral sizes with monolayer flakes observed as well as thicker flakes. SEM analysis, seen in the Supporting Information (Figure S3), of the membrane formed after filtration shows an  $MoS<sub>2</sub>$  "paper" made up of [many small](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)



**Figure 2.** High resolution XPS spectra of MoS<sub>2</sub> raw powder showing the S 2p (a) and Mo 3d (b) binding regions. The suborbital splitting is labeled. The same binding energy regions are shown for the solution exfoliated material in (c) and (d).

individual flakes that are stacked horizontally on top of one another. The individual lateral flake size can be seen to range from 200 nm up to 1  $\mu$ m; this is in agreement with literature values for flakes prepared under similar conditions.<sup>19,42</sup>

To investigate the chemical state of the  $MoS<sub>2</sub>$  films, XPS was performed on the raw powder and the exfoli[ated](#page-9-0) flakes. Previously, it has been shown that XPS can be used to determine the coordination of the molybdenum atoms to distinguish between the different phases that exist, either trigonal prismatic in the case of  $2H-MoS<sub>2</sub>$  or octahedral for 1T- $\text{MoS}_{2}$ ,  $^{12,43}$  XPS can also be used to determine the purity of the starting material, which can greatly affect the subsequent perfor[man](#page-9-0)ce in electrochemical energy storage devices, e.g. presence of redox active contaminants. First, wide-scale survey spectra were taken over a large range of binding energies to determine the presence of different elements (Supporting Information Figure S4). As expected, the only elements present were molybdenum, sulfur, as well as adventitious [carbon and](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf) [oxygen. As](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf) well as the wide-scale spectra, high resolution spectra were taken of the main elements; Figure 2 shows the sulfur 2p (S 2p) and the molybdenum 3d (Mo 3d) shells for the raw  $MoS_2$  powder (Figure 2a,b) and the film after exfoliation (Figure 2c,d). When we compare the spectra for the raw powder and the exfoliated films, we can see a clear difference. In particular the sulfur 2p typically splits into two components in  $MoS_2$  (S 2 $p_{1/2}$  and the S 2 $p_{3/2}$ ); however, in the raw powder four separate peaks are resolved indicating that the sulfur atoms are in two different chemical states. This is likely caused by the partial oxidation of the  $M_0S_2$  into  $M_0S_xO_y$  at edges and defect sites; this is supported by the molybdenum spectra shown in Figure 2b where we see the expected Mo 3d<sub>5/2</sub> and Mo 3d<sub>3/2</sub> as well as a third peak at ∼236 eV which can be attributed to molybdenum  $(Mo^{6+})$  in an octahedral configuration, which is typically observed in partially oxidized  $MoS<sub>2</sub><sup>46–48</sup>$  After exfoliation the XPS spectra of the film change significantly with both the S 2p (Figure 2c) and Mo 3d (Figure 2d) r[egi](#page-9-0)[ons](#page-10-0) representing the expected spectra for pristine 2H- $MoS<sub>2</sub><sup>43,46</sup> We can see in Figure 2d that a small peak still$ remains at ∼236 eV, and this indicates that a very small amount of sur[face](#page-9-0) oxide is still present; however, the intensity is greatly reduced from that of the starting material. This result indicates that the solution phase exfoliation removes any surface oxides that have formed on the  $MoS<sub>2</sub>$  powder as well as removes any amorphous or contaminated material.

To enhance the performance of the semiconducting 2H- $MoS<sub>2</sub>$ , we have produced composites of  $MoS<sub>2</sub>$  and graphene in differing ratios and characterized the resulting membranes by Raman and XPS before electrochemical testing. Composite membranes were synthesized by mixing dispersions of  $M_0S_2$ and graphene in differing weight ratios: 1:1, 1:3, 3:1  $(MoS<sub>2</sub>/$ graphene). The performance of each of these composite membranes was then compared to the respective "pure" membranes (data not shown). It was found that each of the composites had improved performance over the pure membranes, but the 1:3 composite membranes exhibited the highest specific capacitance, and so we will focus on this ratio. The superior performance of the 1:3 ratio composite is attributed to the balance between the high concentration of highly conductive graphene combined with a smaller amount of  $MoS<sub>2</sub>$  flakes to prevent restacking and provide extra surface area. Commercially available graphene powder was utilized and first characterized by Raman and XPS before use, as discussed further in the Supporting Information (Figure S5). The graphene was found to consist of flakes that are approximately

<span id="page-4-0"></span>

Figure 3. (a) Optical image ( $\times$ 100 magnification) of composite membrane ( $\textrm{MoS}_2/\textrm{graph}$ ene ratio, 1:3) supported on PVDF filtration membrane. Raman maps showing distribution of  $MoS<sub>2</sub> A<sub>1g</sub>$  intensity (b) and graphene G band intensity (c) within the composite. (d) Raman spectrum of composite membrane showing both the MoS<sub>2</sub> peaks (enlarged in inset, E<sup>1</sup><sub>2g</sub>, 381 cm<sup>−1</sup>; A<sub>1g</sub>, 405 cm<sup>−1</sup>) and graphene peaks (D, 1348 cm<sup>−1</sup>; G, 1582 cm<sup>-1</sup>; 2D, 2700 cm<sup>-1</sup>).

3−5  $\mu$ m in length, with thicknesses of a few layers (1–3). XPS analysis indicated that some residual contaminants from the exfoliation process remained; however, these were removed after dispersing into NMP and depositing onto the substrate. Membranes of the graphene were created in the same fashion as for the  $MoS<sub>2</sub>$  and composites, and each was characterized separately.

To determine the distribution of  $MoS<sub>2</sub>$  and graphene within the composite membranes we can use Raman spectroscopy to take a map over a large area. These confocal Raman maps are created by taking spectra every  $0.5 \ \mu m$  in a raster pattern over a selected area. The intensity of a given peak can then be plotted visually, and used to determine the spatial distribution within the membrane. Figure 3a shows the optical photograph of one of the composite membranes, using a 100x objective. The high surface roughness seen is ideal for thin layer supercapacitor devices, although some highly reflective larger restacked flakes are visible and the performance of the membrane could be further improved by minimizing these large aggregates. The area shown by the square was then mapped and the corresponding maps are shown in Figure 3b ( $MoS<sub>2</sub>$ ,  $A<sub>1g</sub>$ intensity) and Figure 3c (graphene, G band intensity). From these maps we can see that both the  $MoS<sub>2</sub>$  and graphene are homogeneously distributed throughout the membrane, even the darker areas of the maps represent only a minimum in the signal not an absence of either material. At all areas measured there was a combination of both graphene and  $MoS<sub>2</sub>$  signal. This indicates that during the mixing of the each of the dispersions and subsequent membrane formation, the two materials are thoroughly mixed and evenly distributed. Figure 3d shows a representative Raman spectrum taken from the composite membrane, showing both the characteristic  $MoS<sub>2</sub>$ peaks and the graphene peaks as labeled.

To further investigate the morphology of each membrane, SEM analysis was performed, and typical SEM images of each membrane composition are shown in Figure 4. Figure 4a shows the graphene membrane, and we can clearly see that it consists of many large sheets  $(3-5 \mu m)$  forming a highly wrinkled



Figure 4. SEM images of the graphene membrane (a), MoS2 membrane  $(b)$ , and the composite  $(1:3)$  membrane  $(c)$ . The scale bar in each image is 5  $\mu$ m.

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Figure 5. (a) Photograph of assembled symmetrical coin cell (CR2032). Exploded view schematically shows the internal construction. Stainless steel spacers are inserted to ensure good electrical contact directly between the active material and the outer electrode. The PVDF membranes themselves act as the ion transporting spacers. (b, c) Photographs of PVDF supported membrane formed from exfoliated MoS<sub>2</sub> (mass loading of 1 mg/cm<sup>2</sup>, , scale bar is 13 mm). (d) Plot comparing the cyclic voltammograms (CVs) of the symmetrical coin cells constructed from MoS<sub>2</sub> (blue), graphene (black), and composite membranes (1:3) showing the increased capacitance of the composite membrane. Scan rate is 50 mV/s, and the electrolyte is aqueous 1 M  $Na<sub>2</sub>SO<sub>4</sub>$ .



Figure 6. Plot of specific capacitance change (%) with continued charge/discharge cycles (current density of 1 mA/cm<sup>2</sup>) for the symmetrical MoS<sub>2</sub> cell (a), graphene cell (b), and composite (1:3) membrane cell (c). In each case there is a large increase in specific capacitance with continued cycling before stabilization. The CVs (scan rate of 50 mV/s) precycle and post 10 000 cycles are shown for MoS<sub>2</sub> (d), graphene (e), and the composite (f).

surface. Figure 4b shows the  $MoS<sub>2</sub>$  membrane, and because of the smaller flake size (0.2−1  $\mu$ m), this forms a smoother continuous film with several cracks formed due to manipulation during t[he](#page-4-0) [samp](#page-4-0)le preparation. The composite film is seen in Figure 4c, and because of the combination of the larger graphene sheets and the smaller  $MoS<sub>2</sub>$  flakes, we observe a [rougher s](#page-4-0)urface with increased surface area compared to either of the "pure" membranes. Smaller individual  $MoS<sub>2</sub>$  flakes are also visible stacked on top of the larger graphene sheets. From the high resolution images shown in the Supporting Information (Figure S3) we can also see that the thickness of these films is on average  $\sim$ 0.5  $\mu$ m.

[Electroch](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)emistry. After thorough characteriz[ation](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf) [of](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf) [the](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)  $MoS<sub>2</sub>$  suspensions and resulting films, electrochemical testing was carried out in a symmetrical coin cell architecture (CR2032), as seen in Figure 5a. Individual electrodes were produced by filtering the suspension of exfoliated material

through a preweighed polyvinylidene fluoride (PVDF) filter to form a  $MoS<sub>2</sub>$  membrane, without any additional polymeric binders or conductive additives. Approximately 0.75 mg (mass density of 1 mg/cm<sup>2</sup>) of exfoliated  $\mathrm{MoS}_{2}$  was used as the active material (an optical image of one of the  $MoS<sub>2</sub>$  membranes can be seen in Figure 5b). These flexible membranes were then stacked back-to-back within the cell, with the supporting membranes themselves acting as the ion porous charge separator and providing direct electrical contact between the current collector and the active material. The electrolyte used in this work was aqueous 1 M  $Na<sub>2</sub>SO<sub>4</sub>$ ; a pH neutral electrolyte was used to avoid side reactions such as hydrogen evolution, increasing the long-term stability of the membrane. However, because of the relatively large size of the solvated electrolyte ions, when compared to strong acidic or basic electrolytes, the specific capacitance of  $Na<sub>2</sub>SO<sub>4</sub>$  can be more than 10 times lower than for  $H_2SO_4$ .<sup>30</sup> To act as a comparison, the

performance of a  $MoS_2$  coin cell with 1 M  $H_2SO_4$  can be seen in the Supporting Information (Figure S6); the current density is seen to be ∼10 times higher than the neutral electrolyte. Howe[ver, hydrogen evolution](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf) occurs when the potential is increased, as  ${\rm MoS}_2$  has been shown to be an excellent  ${\rm H}_2$ evolution catalyst.<sup>9,12,49</sup> This  $H_2$  evolution severely limits the lifetime of the coin cell because of deformation of the cell with increasing gas [pr](#page-8-0)[es](#page-9-0)[su](#page-10-0)re. It has also been demonstrated previously that the potential window can be extended up to 1.6 V with the use of sodium sulfate  $(Na_2SO_4)$  as the electrolyte, because of the higher overpotential for the electrolysis of water.<sup>50</sup> Organic electrolytes allow for a large potential window to be used when testing supercapacitor electrodes, as they a[re](#page-10-0) not limited by the electrolysis of water which limits the operating potential of aqueous devices to ∼1 V.<sup>6,24</sup> However, organic electrolytes are hampered by low conductivity and bulky ions, as well as sensitivity to c[on](#page-8-0)[ta](#page-9-0)mination from atmospheric moisture.

Figure 5d compares the cyclic voltammograms (CV) of the  $MoS<sub>2</sub>$  membranes as well as a graphene film and a composite of b[oth graph](#page-5-0)ene and  $MoS<sub>2</sub>$ . The response of the cell shows a typical "square" EDLC response for each of the membranes. When the three membranes are compared, the composite shows a significantly improved performance over either of the individual components. One key drawback to pure  $MoS<sub>2</sub>$  based supercapacitors is the inherent resistivity of the films when compared to carbon based devices. Conductivity of the film can be increased by mixing of graphene and  $MoS<sub>2</sub>$  together to form composites while also exploiting the inherent capacitance of graphene. Several examples of such composites have been demonstrated previously, where it is shown that the addition of conductive material increases the specific capacitance.<sup>36,38,51</sup>

Cycling Stability. One of the most important metrics for comparing supercapacitor electrodes is the cycling [stab](#page-9-0)[ili](#page-10-0)ty. Ideal EDLC supercapacitors should possess extremely high stability with minimal degradation in performance with repeated charge/discharge cycles, while pseudocapacitors typically sacrifice cycling stability for increased energy density.<sup>4</sup> Typically when graphene or  $MoS<sub>2</sub>$  supercapacitor electrodes are cycled continuously, there is a gradual degradation i[n](#page-8-0) capacitance due to increased resistance between flakes, collapse of the porous structure, as well as desorption of active material from the current collector. Thus, the cycling stability is a good indication of the material's suitability for real world supercapacitor applications. Previously, reported literature values for graphene supercapacitor cycling stability vary depending on the testing conditions but can range from 30% loss to almost no loss after 10 000 cycles.<sup>1</sup> This is compared to a 30% decrease in capacitance for solution exfoliated  $MoS<sub>2</sub>$  after only 200 cycles<sup>35</sup> and an 8% decrease [o](#page-8-0)ver 1000 cycles for laser patterned  $MoS<sub>2</sub>.<sup>31</sup>$ 

Figure 6 shows the cycling stability of the  $MoS<sub>2</sub>$  (Figure 6a), graph[ene](#page-9-0) (Figure 6b), and composite (Figure 6c) membranes o[ver 10 00](#page-5-0)0 charge/discharge cycles at a current d[ensity of](#page-5-0) 1 mA/cm<sup>2</sup>. [The interm](#page-5-0)ediate discharge cy[cles \(ever](#page-5-0)y 2000 cycles) are shown for each membrane in the Supporting Information (Figure S10). In the case of the  $MoS<sub>2</sub>$  membrane (Figure 6a) instead of a gradual degradation of the specifi[c capacitance with](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf) repeated cycling as seen typically for such devices, 35,36,52 there is an initial large increase (800%) in the performan[ce](#page-5-0) [over](#page-5-0) [th](#page-5-0)e first ∼3000 cycles before stabilizing for the re[maini](#page-9-0)[ng](#page-10-0) 7000 cycles. Upon assembly electrochemical cells require "preconditioning" (several charge/discharge cycles) to ensure the

electrolyte has permeated throughout the porous structure and the cell has stabilized. $^{24}$  For all results presented here each system was preconditioned by charging/discharging for ∼200 cycles  $(1 \text{ mA/cm}^2)$  prior [to](#page-9-0) further characterization. Thus, the observed increase with repeated cycling only occurs when >1000 charge/discharge cycles are applied. An approximate doubling of the specific capacitance  $(10 \text{ mF/cm}^2)$  is measured for the graphene membrane (Figure 6b) over the first ∼500 cycles, after which the performance increases only slightly over the remaining cycles. The c[omposite](#page-5-0) membrane exhibits a different behavior, increasing gradually over the duration to  $\sim$ 250% of the initial value ( $11 \text{ mF/cm}^2$ ) after 10 000 cycles. The lack of any decrease with continued cycling demonstrates the applicability of these membranes to devices.

The large initial increase in capacitance of the  $MoS<sub>2</sub>$ membrane can be attributed to a combination of (1) a transition in the charge storage mechanism from EDLC to pseudocapacitive ion intercalation and deintercalation, combined with (2) increased active surface area due to partial exfoliation. Initially the predominant charge storage mechanism of the  $MoS<sub>2</sub>$  membrane is accumulation of ions at the doublelayer interface between the flakes and the electrolyte, as indicated by the shape of the CV plot (seen in Figure 5d). However, during charging alkali metal ions present in the electrolyte  $(Na^+)$  can adsorb onto the surface an[d intercala](#page-5-0)te between the  $MoS<sub>2</sub>$  layers of the membrane, represented by eq 1.

$$
MoS_2 + Na^+ + e^- \leftrightarrow MoS-SNa
$$
 (1)

This ion intercalation pseudocapacitance has been observed previously for  $\text{MoS}_{2}$  films<sup>29−31,37</sup> and particularly for transition n<br>metal oxide films.<sup>22,53–55</sup> The continual intercalation and deintercalation of the rela[tively la](#page-9-0)rge sodium ions over repeated cycles can also lead [to](#page-9-0) [exfoli](#page-10-0)ation of the layers that are restacked during the filtration process,<sup>56,57</sup> leading to an increased surface area and subsequently a much larger specific capacitance. Layered materials such as  $MoS<sub>2</sub>$  are unique in that their structure allows for this dynamic expansion to allow ions to intercalate without damage to the lattice occurring.<sup>29</sup> This transition in charge storage mechanism can be seen when comparing the CV plots before and after cycling, as [se](#page-9-0)en in Figure 6d, which show that after cycling the current has increased significantly, but the shape of the plot is less ideal, [indicative](#page-5-0) of pseudocapacitive Faradaic charge transfer.  $21,22$ Further evidence of this partial exfoliation can be seen in the SEM analysis shown in the Supporting Information (Fi[gure](#page-9-0) S11), where we see that after cycling the  $MoS<sub>2</sub>$  membrane exhibits increased surface rou[ghness. The graphene me](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)mbrane also shows an increase in capacitance (Figure 6b) after cycling, and more noticeable is the change in shape of the CV plot (Figure 6e). Because of its high cond[uctivity, th](#page-5-0)e precycle CV plot displays highly ideal double-layer charge storage, as [expected](#page-5-0) from graphene.<sup>4,58</sup> After cycling there is an increase in capacitance of ∼250%; again the shape of the CV curve after cycling indicates pseud[o](#page-8-0)[cap](#page-10-0)acitive behavior. This increased capacitance with continuous cycling has been reported previously for reduced graphene oxide films and was also attributed to the partial exfoliation with continued ion intercalation/deintercalation increasing the active surface area, termed "electroactivation". 33,59,60

Finally, the composite membrane (Figure 6c) shows a moderate steady increase [wi](#page-9-0)[th cy](#page-10-0)cling duration up to ∼250%. When we compare the CV data before [and after](#page-5-0) the cycling

(Figure 6f), we see that the response indicates highly ideal EDLC behavior when compared to the pure membranes in [Figure 6d](#page-5-0) and Figure 6e. This indicates that the composite of both graphene and  $MoS<sub>2</sub>$  together is better able to resist [structura](#page-5-0)l cha[nges tha](#page-5-0)t occur with continual charge and discharge cycles. Typically, when dispersions of 2D materials are deposited into a film, the van der Waals attraction between neighboring sheets causes restacking to occur, and this limits the available surface area, thus reducing the specific capacitance. Thus, it is deduced that differences in flake dimensions and surface energy in the case of the  $MoS<sub>2</sub>/graph$ ene composites form disordered heterostructures and prevent restacking, as seen previously in the SEM analysis (Figure 4). It is this physical prevention of restacking, effectively increasing the specific surface area, that is likely resp[onsible fo](#page-4-0)r the ideal EDLC charge storage behavior observed.

Specific Capacitance of Thin-Layer Membranes. The specific capacitance  $(C_{\text{sp}})$  of the membranes can be calculated from the CV plots using eq  $2<sup>1</sup>$ 

$$
C_{\rm sp} = \frac{\int I(V) \, \mathrm{d}v}{v A \, \Delta V} \tag{2}
$$

where  $\int I(V)$  is equal to the integrated area under the CV plot,  $\nu$  is the scan rate, A is the electrode area, and  $\Delta V$  is the potential window. Figure 7 plots the calculated specific



Figure 7. Plot of specific capacitance  $(mF/cm<sup>2</sup>)$  with increasing scan rate calculated from cyclic voltammetry for each cell before and after cycling 10 000 times (seen in Figure 6).

capacitance with increasin[g scan r](#page-5-0)ates for each of the three membranes before and after cycling. (The CV plots showing the increasing potential scan rates are shown in the Supporting Information, Figure S7−S9.) As expected from eq 2, there is a decrease in specific capacitance with increasing scan [rate, due to](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf) [the transpo](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)rt limit on the supply of ions. As discussed previously, the graphene displays ideal EDLC behavior before cycling, and this is also supported by the high discharge rate as the scan rate increases, with the specific capacitance decreasing by only 43% when increasing the scan rate from 5 mV/s (1.16 mF/cm<sup>2</sup>) up to 1000 mV/s (0.66 mF/cm<sup>2</sup>). The  ${\rm MoS}_2$ membrane, however, displays very high specific capacitance at low scan rates  $(1.83 \text{ mF/cm}^2 \text{ at } 5 \text{ mV/s})$  but decreases by 80% at 1000 mV/s to 0.37 mF/cm<sup>2</sup>, which is indicative of pseudocapacitive behavior. $^{30}$  These capacitance values correlate well with reported values for solution exfoliated  $MoS<sub>2</sub>$  films that were able to produce  $\sim$ 2 [mF](#page-9-0)/cm<sup>2</sup> at 10 mV/s.<sup>35</sup> The composite membrane exhibits much larger specific capacitance than either of the single component membranes; the s[pec](#page-9-0)ific capacitance was measured as  $4.29 \text{ mF/cm}^2$  at 5 mV/s; however, the

composite suffers from the same degradation as seen in the pure  $MoS<sub>2</sub>$  membrane with increasing scan rate, decreasing by ∼80% at 1000 mV/s to 0.98 mF/cm<sup>2</sup> . This larger specific capacitance precycling is attributed to physical interactions between the  $MoS<sub>2</sub>$  and graphene flake, preventing restacking from occurring and increasing the active surface area. As well as the area normalized specific capacitances, values for the gravimetric and volumetric capacitance are provided in the Supporting Information (Table S1−S2) to allow for comparison to other similar works.

[The postcycle speci](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)fic capacitance increases significantly for all three membrane systems, attributed to "electroactivation" as discussed previously, and the performance of the  $MoS<sub>2</sub>$  and composite membranes are seen to match one another closely. This is surprising, as the composite membrane consists of just 25 wt % of  $MoS<sub>2</sub>$  compared to the pure membrane but is able to match its performance. The specific capacitance of the composite membrane ranges from 11 mF/cm<sup>2</sup> (5 mV/s) down to 3.1 mF/cm<sup>2</sup> (1000 mV/s). These specific capacitance values compare favorably with values published previously, ranging from 2 to 3.1 mF/cm<sup>2</sup> for solution exfoliated  $MoS<sub>2</sub><sup>31,35</sup>$  The specific capacitance of the composite also compares favorably to graphene/nanotube composites, which generated [2.16](#page-9-0) mF/  $\text{cm}^{2,32}$  As well as cyclic voltammetry, galvanostatic charge/ . discharge at different current densities was used to calculate the speci[fi](#page-9-0)c capacitance of each of the membranes (see Supporting Information, Figure S7−S9), and the values were seen to correspond closely to those calculated from C[V. Electro](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)[chemical im](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)pedance spectroscopy (EIS) was also used to investigate the frequency response of the coin cells before and after cycling (Supporting Information, Figure S12−S13), supporting the transition in charge storage mechanism from EDLC to pseud[ocapacitance. From the EI](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b04672/suppl_file/am5b04672_si_001.pdf)S analysis we can also determine the equivalent series resistance (ESR) within the cells, with the  $MoS<sub>2</sub>$  membrane having the highest ESR of 1.6  $Ω$  and lower for the composite (0.7  $Ω$ ) and pure graphene  $(0.96 \Omega)$ .

Self-Discharge Behavior. Although very important when discussing application of supercapacitor designs, the selfdischarge behavior is often neglected when characterizing the performance of a given material. $4,25$  Self-discharging of supercapacitors can severely limit their applicability to real world applications and thus are an i[m](#page-8-0)[po](#page-9-0)rtant figure of merit when comparing electrode designs. The spontaneous voltage decrease that occurs within a charged supercapacitor under open circuit conditions, often referred to as self-discharge, can occur because of three mechanisms: Faradaic processes, leakage current, and charge redistribution. <sup>61,62</sup> These Faradaic processes that can occur, when aqueous electrolyte solutions are used, are thought to be predomina[ntly o](#page-10-0)xygen reduction at the negative electrode. $63$  Leakage current occurs when charged ions or impurities within the electrolyte spontaneously migrate from each electrode, [th](#page-10-0)ereby reducing the cell potential. $62$ Charge redistribution occurs when selective areas of the active material are unevenly charged, leading to areas that [are](#page-10-0) overcharged while others have lower charge. This charge redistribution between these areas then leads to an overall decrease in cell potential. Thus, when a supercapacitor electrode material is characterized, minimizing the selfdischarge rate is important to maximize performance. In the case of EDLC devices the voltage decreases initially from its initial value quickly and the rate decreases with time.<sup>61</sup> Over this initial rapid decrease in potential the device is exhibiting

<span id="page-8-0"></span>Figure 8 plots the open-circuit voltage of each of the three symmetrical cells (postcycling); each underwent 15 min of



Figure 8. Self-discharge behavior for the symmetrical graphene cell (black),  $MoS<sub>2</sub>$  cell (blue), and composite cell (red). The cells each underwent potentiostatic charging at 1 V for 15 min prior to measuring the open-circuit potential.

potentiostatic charging to 1 V before the open-circuit potential was measured over 1 h. It can be seen that the graphene membrane (black) suffers from the highest rate of selfdischarge, followed closely by the  $MoS<sub>2</sub>$  membrane (blue). The composite membrane (red), however, displays much better charge retention (37% greater than the graphene, 25% greater than  $MoS<sub>2</sub>$ ). The increased charge retention performance of the composite membrane indicates that the synergistic effect of combing the highly conductive graphene with the pseudocapacitive  $MoS<sub>2</sub>$  allows for reduced Faradaic leakage and charge redistribution across the membrane.

# ■ CONCLUSION

We have demonstrated that composite membranes made from solution exfoliated dispersions of 2D materials, such as semiconducting  $MoS<sub>2</sub>$  and semimetallic graphene, can create a synergistic effect that allows for high performance electrochemical energy storage devices to be constructed. The high specific capacitance of the  $MoS<sub>2</sub>/graph$ ene membranes, with only a small mass loading (∼1 mg/cm<sup>2</sup>), no added binders, and no extra porous charge separators, demonstrates a cheap, simple, and scalable electrochemical energy storage cell architecture with excellent cycle life. It was found that the optimal ratio of the  $MoS<sub>2</sub>$  and graphene, in terms of electrochemical performance, was 1:3 ( $MoS<sub>2</sub>/graphene$ ). The composite membrane was able to achieve a high specific capacitance ( $\sim$ 11 mF/cm<sup>2</sup> at 5 mV/s). With continued cycling it was observed that electroactivation due to ion intercalation leads to a significant increase in specific capacitance. The selfdischarge behavior also demonstrated that the composite membrane possessed superior charge retention compared to the individual components. This work suggests that composites of various 2D materials (TMDs, graphene, hBN, etc.) can easily

be combined and used to enhance the performance of single component materials for electrochemical energy storage.

#### ■ ASSOCIATED CONTENT

# **6** Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsami.5b04672.

Raman comparison of powder and single-crystal  $MoS<sub>2</sub>$ fl[akes, Raman map](http://pubs.acs.org)s of co[mposite membranes, furth](http://pubs.acs.org/doi/abs/10.1021/acsami.5b04672)er SEM analysis, comparison of electrolytes, Raman and XPS characterization of graphene powder, galvanostatic charge/discharge curves, and electrochemical impedance spectroscopy for coin cell devices (PDF)

# ■ AUTHOR INFORMATION

#### Corresponding Authors

\*M.A.B.: e-mail, Mark.Bissett@manchester.ac.uk.

\*R.A.W.D.: e-mail, Robert.Dryfe@manchester.ac.uk.

### Author Contrib[utions](mailto:Mark.Bissett@manchester.ac.uk)

All authors have g[iven approval to the](mailto:Robert.Dryfe@manchester.ac.uk) final version of the manuscript.

# Notes

The authors declare no competing financial interest.

# ■ ACKNOWLEDGMENTS

The authors acknowledge funding from BGT Materials, U.K., as well as the EPSRC (Grant ref EP/K016954/1). The authors also thank Anders J. Barlow at the National EPSRC XPS User's Service (NEXUS) at Newcastle University for assistance acquiring the XPS data.

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